

Hyperkähler reduction and the moduli of flat bundles on complex curves

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Plan

1. Symplectic reduction and GIT
2. Hyperkähler reduction
3. Moment map for the gauge group action.

Conventions: Further on, G is a compact, connected Lie group, $G_{\mathbb{C}}$ its complexification, \mathfrak{g} and $\mathfrak{g}_{\mathbb{C}}$ the corresponding Lie algebras. **Central element** of \mathfrak{g}^* is one which is fixed by the adjoint action of G .

Cartan's formula and symplectomorphisms

We denote the Lie derivative along a vector field as $\text{Lie}_x : \Lambda^i M \longrightarrow \Lambda^i M$, and contraction with a vector field by $i_x : \Lambda^i M \longrightarrow \Lambda^{i-1} M$.

Cartan's formula: $d \circ i_x + i_x \circ d = \text{Lie}_x$.

REMARK: Let (M, ω) be a symplectic manifold, G a Lie group acting on M by symplectomorphisms, and \mathfrak{g} its Lie algebra. For any $g \in \mathfrak{g}$, denote by ρ_g the corresponding vector field. Then $\text{Lie}_{\rho_g} \omega = 0$, giving $d(i_{\rho_g}(\omega)) = 0$. **We obtain that $i_{\rho_g}(\omega)$ is closed, for any $g \in \mathfrak{g}$.**

DEFINITION: **A Hamiltonian** of $g \in \mathfrak{g}$ is a function h on M such that $dh = i_{\rho_g}(\omega)$.

Moment maps

DEFINITION: (M, ω) be a symplectic manifold, G a Lie group acting on M by symplectomorphisms. **A moment map** μ of this action is a linear map $\mathfrak{g} \rightarrow C^\infty M$ associating to each $g \in \mathfrak{g}$ its Hamiltonian.

REMARK: It is more convenient to consider μ as an element of $\mathfrak{g}^* \otimes_{\mathbb{R}} C^\infty M$, or (and this is most standard) **as a function with values in \mathfrak{g}^*** .

REMARK: Moment map **always exists** if M is simply connected.

DEFINITION: A moment map $M \rightarrow \mathfrak{g}^*$ is called **equivariant** if it is equivariant with respect to the coadjoint action of G on \mathfrak{g}^* .

REMARK: $M \xrightarrow{\mu} \mathfrak{g}^*$ is a moment map iff for all $g \in \mathfrak{g}$, $\langle d\mu, g \rangle = i_{\rho_g}(\omega)$. Therefore, **a moment map is defined up to a constant \mathfrak{g}^* -valued function**. An equivariant moment map is defined up to **a constant \mathfrak{g}^* -valued function which is G -invariant**, that is, up to addition of a central vector $c \in \mathfrak{g}^*$.

CLAIM: **An equivariant moment map exists whenever $H^1(G, \mathfrak{g}^*) = 0$** . In particular, when G is reductive and M is simply connected, an equivariant moment map exists. Further on, all moment maps will be tacitly considered equivariant.

Weinstein-Marsden theorem

DEFINITION: (Weinstein-Marsden) (M, ω) be a symplectic manifold, G a compact Lie group acting on M by symplectomorphisms, $M \xrightarrow{\mu} \mathfrak{g}^*$ an equivariant moment map, and $c \in \mathfrak{g}^*$ a central element. The quotient $\mu^{-1}(c)/G$ is called **symplectic reduction** of M , denoted by $M//G$.

CLAIM: The symplectic quotient $M//G$ is a symplectic manifold of dimension $\dim M - 2 \dim G$.

Proof. Step 1: $T_x(\mu^{-1}(c)) = d\mu^{-1}(0)$, however, $d\mu$ is ω -dual to the space $\tau(\mathfrak{g})$ of vector fields tangent to the G -action, **hence** $d\mu^{-1}(0) = \tau(\mathfrak{g})^\perp$.

Step 2: Since μ is G -equivariant, G preserves $\mu^{-1}(c)$, hence $\tau(\mathfrak{g}) \subset d\mu^{-1}(0)$. This implies that $\tau(\mathfrak{g}) \subset TM$ **is isotropic** (that is, $\omega|_{\tau(\mathfrak{g})} = 0$.) Its ω -orthogonal complement in $T_x M$ is $T_x(\mu^{-1}(c))$ (Step 1).

Step 3: Consider the **characteristic foliation** \mathcal{F} on $\mu^{-1}(c)$, that is, the set of all $v \in T_x(\mu^{-1}(c))$ such that $\omega(v, w) = 0$ for all $w \in T_x(\mu^{-1}(c))$. From Step 2 **we obtain that** $\mathcal{F} = \tau(\mathfrak{g})$.

Step 4: Since $\omega|_{\mu^{-1}(c)}$ is closed, it satisfies $\text{Lie}_v(\omega) = 0$ for all $v \in \mathcal{F}$. This implies that it is lifted from the leaf space of characteristic foliation, identified with $M//G$. ■

Symplectic reduction and GIT

THEOREM: Let (M, I, ω) be a Kähler manifold, $G_{\mathbb{C}}$ a complex reductive Lie group acting on M by holomorphic automorphisms, and G its compact form acting isometrically. **Then $M//G$ is a Kähler manifold.**

Proof: Since the orbits of the $G_{\mathbb{C}}$ -action are complex subvarieties, they are symplectic. Since the orbits of $G \subset G_{\mathbb{C}}$ are isotropic, and their dimension is half of dimension of orbits of $G_{\mathbb{C}}$, they are actually Lagrangian subvarieties in orbits of $G_{\mathbb{C}}$. Therefore, $\mu^{-1}(c)$ intersects each orbit of $G_{\mathbb{C}}$ in a G -orbit. **We have identified $M//G$ with a space of $G_{\mathbb{C}}$ -orbits which intersect $\mu^{-1}(c)$.**

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REMARK: In such a situation, $M//G$ is called **the Kähler quotient**, or **GIT quotient**. The choice of a central element $c \in \mathfrak{g}^*$ is known as a choice of **stability data**.

REMARK: **The points of $M//G$ are in bijective correspondence with the orbits of $G_{\mathbb{C}}$ which intersect $\mu^{-1}(c)$.** Such orbits are called **polystable**, and the intersection of a $G_{\mathbb{C}}$ -orbit with $\mu^{-1}(c)$ is a G -orbit.

Kähler reduction and a Kähler potential

DEFINITION: Kähler potential on a Kähler manifold (M, ω) is a function ψ such that $dd^c\psi = \omega$.

PROPOSITION: Let G be a real Lie group acting on a Kähler manifold M by holomorphic isometries, and ψ be a G -invariant Kähler potential. **Then the moment map $\mathfrak{g} \times M \xrightarrow{\mu_g} \mathbb{R}$ can be written as $g, m \longrightarrow -\text{Lie}_{Iv} \psi$,** where $v = \tau(g) \in TM$ is the tangent vector field associated with $g \in \mathfrak{g}$.

Proof: Since ψ is G -invariant, and I is G -invariant, we have $0 = \text{Lie}_v d^c\psi = (dd^c\psi) \lrcorner v + d(\langle d^c\psi, v \rangle)$. Using $\omega = dd^c\psi$, we rewrite this equation as $\omega \lrcorner v = -d(\langle d^c\psi, v \rangle)$, giving an equation for the moment map $\mu_g = -\langle d^c\psi, v \rangle$. Acting by I on both sides, we obtain $\mu_g = -\langle d\psi, Iv \rangle = -\text{Lie}_{Iv} \psi$. ■

COROLLARY: Let V be a Hermitian representation of a compact Lie group G . **Then the corresponding moment map can be written as $\mu_g(v) = -\text{Lie}_{Ig} |v|^2 = -\frac{1}{2} \langle v, Ig(v) \rangle$.** ■

Hyperkähler manifolds

DEFINITION: A **hyperkähler structure** on a manifold M is a Riemannian structure g and a triple of complex structures I, J, K , satisfying quaternionic relations $I \circ J = -J \circ I = K$, such that g is Kähler for I, J, K .

REMARK: A hyperkähler manifold **has three symplectic forms**
 $\omega_I := g(I\cdot, \cdot)$, $\omega_J := g(J\cdot, \cdot)$, $\omega_K := g(K\cdot, \cdot)$.

REMARK: This is equivalent to $\nabla I = \nabla J = \nabla K = 0$: the parallel transport along the connection preserves I, J, K .

REMARK:

The form $\Omega := \omega_J + \sqrt{-1} \omega_K$ is **holomorphic and symplectic** on (M, I) .

Hyperkähler reduction

DEFINITION: Let G be a compact Lie group, ρ its action on a hyperkähler manifold M by hyperkähler isometries, and \mathfrak{g}^* a dual space to its Lie algebra. **A hyperkähler moment map** is a G -equivariant smooth map $\mu : M \rightarrow \mathfrak{g}^* \otimes \mathbb{R}^3$ such that $\langle \mu_i(v), g \rangle = \omega_i(v, d\rho(g))$, for every $v \in TM$, $g \in \mathfrak{g}$ and $i = 1, 2, 3$, where ω_i is one three Kähler forms associated with the hyperkähler structure.

DEFINITION: Let ξ_1, ξ_2, ξ_3 be three G -invariant vectors in \mathfrak{g}^* . The quotient manifold $M // G := \mu^{-1}(\xi_1, \xi_2, \xi_3) / G$ is called **the hyperkähler quotient** of M .

THEOREM: (Hitchin, Karlhede, Lindström, Roček)

The quotient $M // G$ is hyperkaehler.

Holomorphic moment map

Let $\Omega := \omega_J + \sqrt{-1}\omega_K$. This is a holomorphic symplectic (2,0)-form on (M, I) .

The proof of HKLR theorem. Step 1: Let μ_J, μ_K be the moment map associated with ω_J, ω_K , and $\mu_{\mathbb{C}} := \mu_J + \sqrt{-1}\mu_K$. Then $\langle d\mu_{\mathbb{C}}, g \rangle = i_{\rho g}(\Omega)$. Therefore, $d\mu_{\mathbb{C}} \in \Lambda^{1,0}(M, I) \otimes \mathfrak{g}^*$.

Step 2: This implies that the map $\mu_{\mathbb{C}}$ is holomorphic. It is called **the holomorphic moment map**.

Step 3: By definition, $M // G = \mu_{\mathbb{C}}^{-1}(c) // G$, where $c \in \mathfrak{g}^* \otimes_{\mathbb{R}} \mathbb{C}$ is a central element. **This is a Kähler manifold**, because it is a Kähler quotient of a Kähler manifold.

Step 4: We obtain 3 complex structures I, J, K on the hyperkähler quotient $M // G$. **They are compatible in the usual way** (an easy exercise). ■

Gauge group

DEFINITION: Let G be a Lie group, and P a principal G -bundle on M . **The gauge group** of P is the group of G -invariant automorphisms of P .

REMARK: Let $G_{\mathbb{C}}$ be a complex Lie group, P a principal $G_{\mathbb{C}}$ -bundle, \mathcal{A} the space of all $G_{\mathbb{C}}$ -invariant connections on P , and $\mathfrak{P}_{\mathbb{C}}$ the bundle of G -invariant vector fields tangent to the fibers of P . **Then \mathcal{A} is a complex affine vector space with linearization $\Lambda^1 M \otimes \mathfrak{P}_{\mathbb{C}}$.** It is equipped with the gauge group action.

DEFINITION: Let P_G be a reduction of the principal $G_{\mathbb{C}}$ -bundle P to G , and $g \rightarrow g^t$ be the corresponding real structure operator on $\mathfrak{P}_{\mathbb{C}}$. Then $\text{Aut}(P_G)$ is called **the real gauge group**, and $\text{Aut}(P)$ **the complex gauge group**. The Killing form on $\mathfrak{P}_{\mathbb{C}}$ is denoted as $a, b \rightarrow \text{Tr}(ab) \in C^\infty M$.

Hyperkähler structure on the space of connections

CLAIM: Let $P_{\mathbb{C}}$ be a complexification of a principal G -bundle P on a compact Riemann surface M , and $\nabla \in \mathcal{A}$ a connection in $P_{\mathbb{C}}$. Denote by $g \rightarrow g^t$ the real involution on $\mathfrak{P}_{\mathbb{C}}$ fixing \mathfrak{P} . **Then the tangent space $T_{\nabla}\mathcal{A}$ is equipped with a real gauge invariant Hermitian form $g(a, b) = \int_M \operatorname{Re} \operatorname{Tr}(a \wedge b^t)$ and a complex linear 2-form $\Omega(u, v) := \int_M \operatorname{Tr}(u \wedge v)$. ■**

DEFINITION: Define quaternionic structure on $T_{\nabla}\mathcal{A}$ as follows. The complex structure I comes from the complex structure on $\mathfrak{P}_{\mathbb{C}}$, and J comes from $\operatorname{Re} \Omega(x, Jy) = g(a, b)$.

REMARK: For $\lambda \otimes a \in T_{\nabla}\mathcal{A} = \Lambda^1(M) \otimes_{\mathbb{C}} \mathfrak{g}_{\mathbb{C}}$, we can write $J(\lambda \otimes a) = I_M(\lambda) \otimes a^t$, where I_M is a complex structure operator on M acting on $\Lambda^1(M)$.

COROLLARY: The manifold \mathcal{A} is equipped with a natural real gauge invariant flat hyperkähler structure. ■

Hyperkähler moment map

REMARK: The tangent space to the gauge group $\text{Aut}(P)$ can be identified with $\mathfrak{P}_{\mathbb{C}}$. Therefore, **the gauge moment map takes values in $\mathfrak{P}_{\mathbb{C}}^* = \mathfrak{P}_{\mathbb{C}}$.**

THEOREM: Let (M, I, ω) be a Riemannian surface equipped with a Hermitian form ω , P a principal $G_{\mathbb{C}}$ -bundle, \mathcal{A} the space of connections on P , and $G \subset G_{\mathbb{C}}$ a compact real form. Consider the hyperkähler structure on \mathcal{A} defined above. **Then the holomorphic moment map associated with the gauge action can be written as $\nabla \rightarrow -\frac{\Theta_{\nabla}}{\omega}$, where $\Theta_{\nabla} \in \Lambda^2 M \otimes \mathfrak{P}_{\mathbb{C}}$ is the curvature of ∇ .**

Proof. Step 1: If $\nabla_1 = \nabla + A$, we have $\theta_{\nabla_1} = \Theta_{\nabla} + \nabla(A) + A \wedge A$. Therefore, differential of Θ_{∇} takes $A \in \Lambda^1 M \otimes \mathfrak{P}_{\mathbb{C}}$ to $\int_M \nabla(A)$. If we pair this to $b \in \mathfrak{P}_{\mathbb{C}}$, we obtain $A \rightarrow \int_M \text{Tr}(b \nabla(A))$. **This is the differential of a moment map evaluated in ∇, b .**

Step 2: The holomorphic symplectic form on \mathcal{A} is expressed as $\Omega(A, B) = \int_M \text{Tr}(A \wedge B)$. For any $b \in \mathfrak{P}_{\mathbb{C}}$, the corresponding vector field on \mathcal{A} is written as $\nabla(b)$, and the corresponding 1-form as $\Omega(\nabla(b), A) = \int_M \text{Tr}(\nabla(b) \wedge A)$. Here $A \in T\mathcal{A}$ is a tangent vector, and $\Omega(\nabla(b), \cdot)$ is considered as a 1-form on \mathcal{A} .

Step 3: It remains to compare the 1-forms obtained in Step 1 and Step 2: $\int_M \text{Tr}(\nabla(b) \wedge A) = -\int_M \text{Tr}(b \nabla(A))$. ■

Space of flat connections

THEOREM: The space of stable flat connections on a compact Riemann surface up to (complex) gauge equivalence **is obtained as $\mathcal{M} = \mathcal{A} // \mathcal{G}$** , where \mathcal{G} is the real gauge group. **In particular, \mathcal{M} is hyperkähler.**

Proof: Denote the space of flat connections by $\mathcal{A}_{\text{fl}} \subset \mathcal{A}$. Flatness is vanishing of the holomorphic moment map, hence $\mathcal{A}_{\text{fl}} // G = \mathcal{A} // \mathcal{G}$. However, $\mathcal{A}_{\text{fl}} // G$ is precisely the space of stable flat connections up to complex gauge equivalence.

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Addendum: Real moment map

PROPOSITION: Since the Hermitian form on \mathcal{A} is fixed by affine transforms, it can be written as dd^c of a quadratic function which is its Kähler potential. The latter can be written as $\psi(\nabla) = \int_M \text{Tr}[\nabla, \nabla^t]$ (here ∇^t is a connection operator which is real conjugate to ∇).

Proof: Second derivative of this quadratic map is $\text{Hess}(x, y) = \int_M \text{Re}(x \wedge y^t)$.

■

PROPOSITION: The real moment map for gauge action on the space of connections is $\frac{[\nabla, \nabla^t]}{\omega}$.

Proof: The gauge Lie algebra action is written as $g(\nabla) = \nabla + \nabla(g)$, for all $g \in \mathfrak{P}_{\mathbb{C}}$. This gives an expression for the real moment map:

$$\mu_g(\nabla) = \frac{d}{dg} \int_M \text{Tr}[\nabla + \sqrt{-1} \nabla(g), \nabla^t - \sqrt{-1} \nabla^t(g)] = \int_M \text{Tr}([\nabla, \nabla^t](g))$$

We obtain that the real moment map $\mu : \mathcal{A} \rightarrow \mathfrak{P}_{\mathbb{C}}$ takes ∇ and puts it to $\frac{[\nabla, \nabla^t]}{\omega}$. ■